



Semiconductor Devices

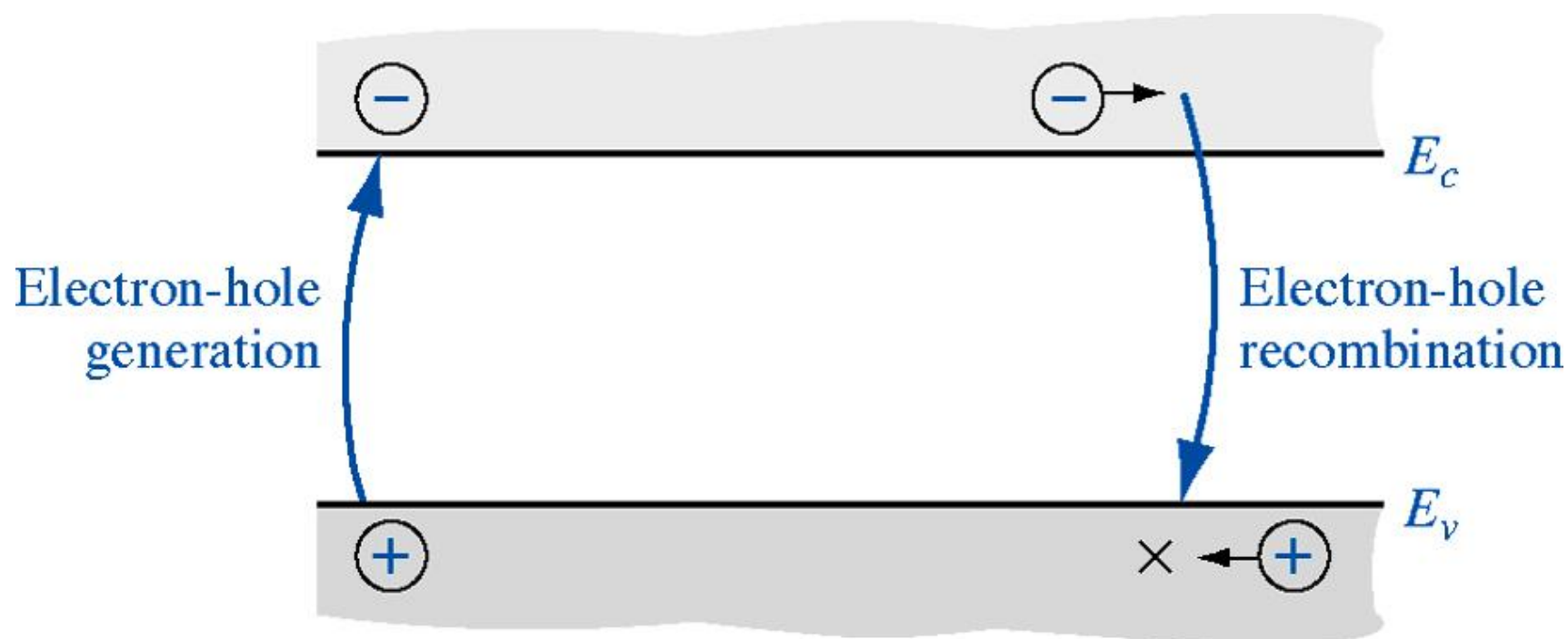
Chapter 6

Nonequilibrium Excess Carriers in Semiconductors

오세용, Ph.D.
한양대 ERICA, 조교수

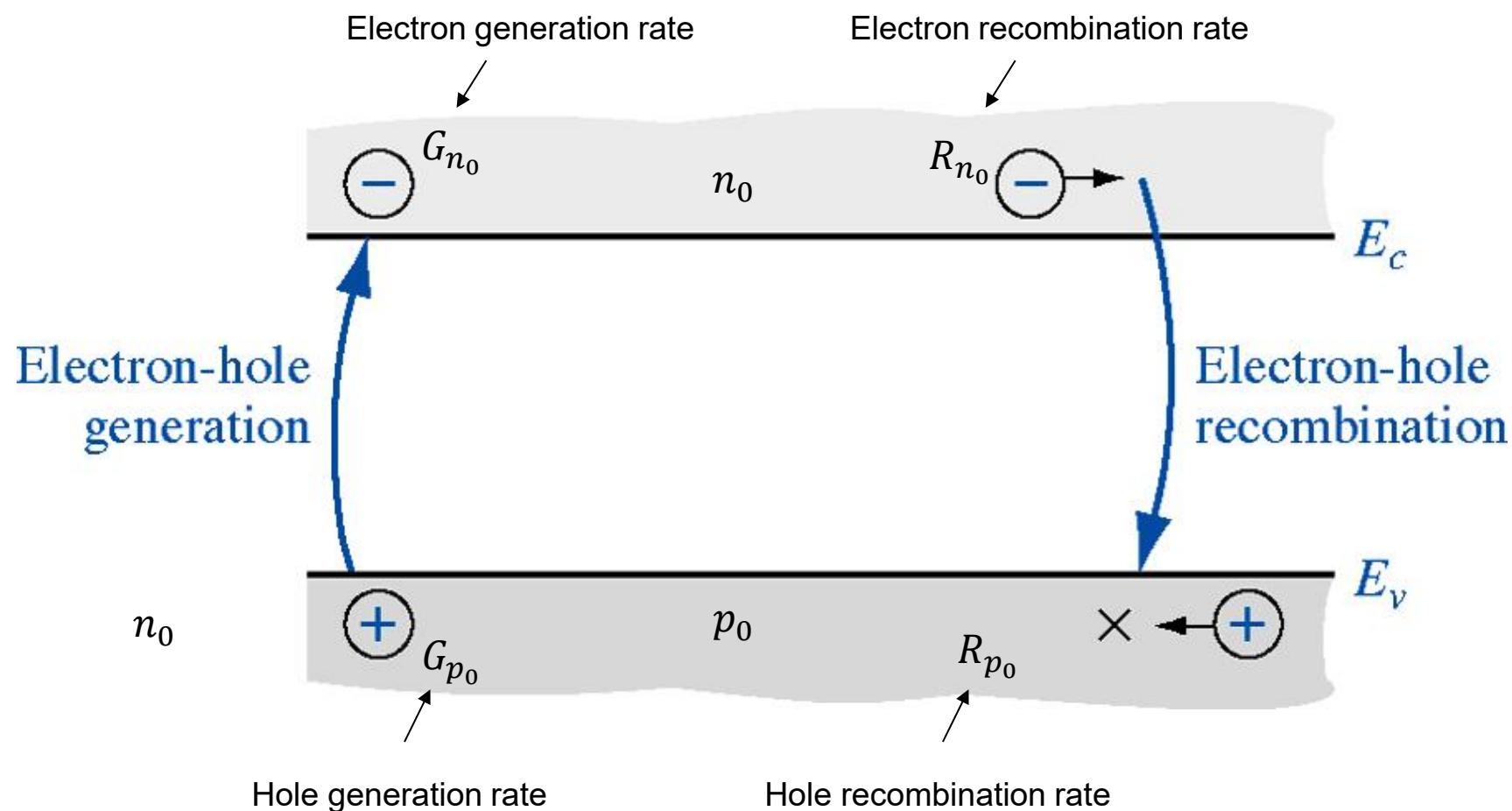


Thermal Equilibrium State



- ✓ **Generation:** produces an Electron-Hole pair
- ✓ **Recombination:** eliminates an Electron-Hole pair

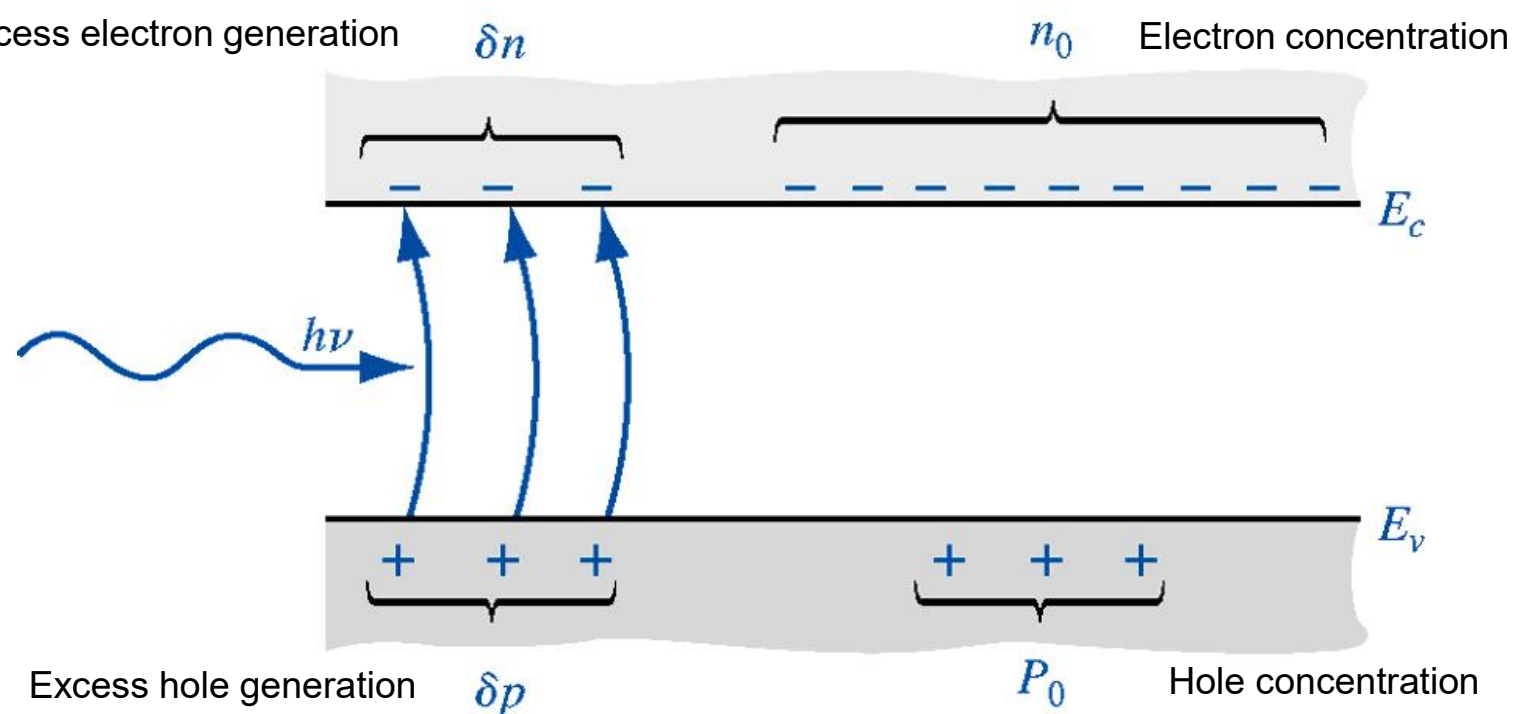
Thermal Equilibrium State



$$G_{n_0} = G_{p_0} = R_{n_0} = R_{p_0} \text{ [}/\text{cm}^3 \cdot \text{s]} \quad (\text{In equilibrium state})$$

$$n_0 p_0 = n_i^2$$

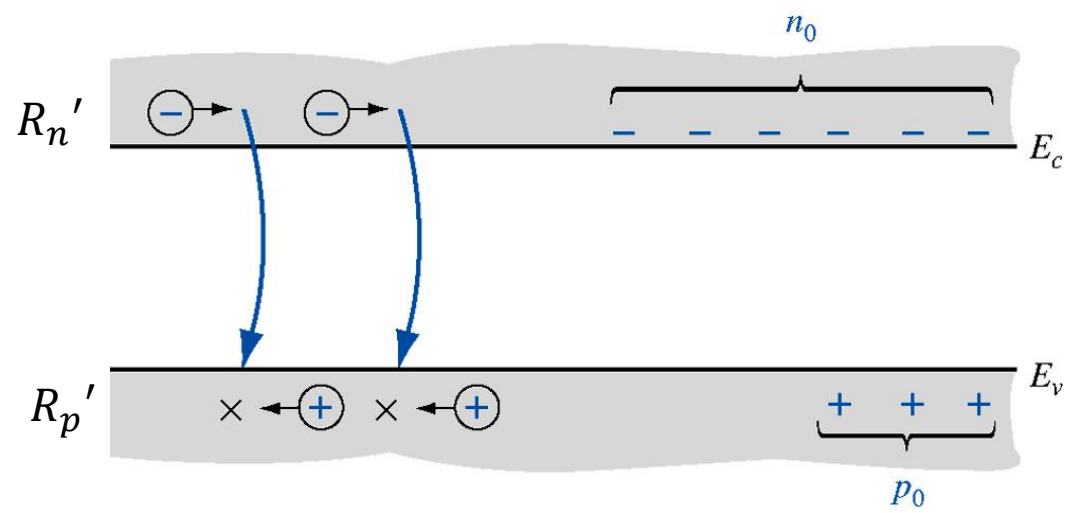
Nonequilibrium State



$$n = n_0 + \delta n, \quad p = p_0 + \delta p \quad \Rightarrow \quad np \neq n_0 p_0 = n_i^2$$

Mass-action law

Recombination Rate of Excess Carriers



$$n = n_0 + \delta n \qquad p = p_0 + \delta p$$



$$np \neq n_0 p_0 = n_i^2$$

Mass-action law

Excess carrier recombination rate

$$R_n' = \frac{\partial(\delta n(t, x))}{\partial t}$$

$$R_p' = \frac{\partial(\delta p(t, x))}{\partial t} \quad [\#/cm^3 \cdot s]$$

↓ simplify

τ_R

Recombination carrier lifetime
→ Mean time for excess carrier recombination

$$R_n' \approx \frac{\delta n}{\tau_R} \qquad R_p' \approx \frac{\delta p}{\tau_R}$$



For n-type Si

For p-type Si

$$R_n' = \frac{\delta n}{\tau_p} = R_p' = \frac{\delta p}{\tau_p}$$

Excess minority carrier lifetime

$$R_n' = \frac{\delta n}{\tau_n} = R_p' = \frac{\delta p}{\tau_n}$$

Generation Rate of Excess Carriers

Excess carrier
generation rate

$$g_n' = \frac{\delta n}{\tau_G}$$

$$g_p' = \frac{\delta p}{\tau_G}$$

[#/cm³ · s]

Generation carrier lifetime

→ Mean time for excess carrier generation

$$g_n' = g_p' = \text{constant} \quad (\text{In general cases})$$

Steady State Condition

Steady-state

$$\frac{dn}{dt} = \frac{dp}{dt} = 0$$

$$\frac{dn}{dt} = (G_{n_0} + g_n') - (R_{n_0} + R_n') = 0 \qquad g_n' = R_n'$$

$$\frac{dp}{dt} = (G_{p_0} + g_p') - (R_{p_0} + R_p') = 0 \qquad g_p' = R_p'$$

Steady-state

→ Regardless of Nonequilibrium or Equilibrium but no net changes in concentrations

Table 6.1 | Relevant notation used in Chapter 6

Symbol	Definition
n_0, p_0	Thermal equilibrium electron and hole concentrations (independent of time and also usually position).
n, p	Total electron and hole concentrations (may be functions of time and/or position).
$\delta n = n - n_0$ $\delta p = p - p_0$	Excess electron and hole concentrations (may be functions of time and/or position).
g'_n, g'_p	Excess electron and hole generation rates.
R'_n, R'_p	Excess electron and hole recombination rates.
τ_{n0}, τ_{p0}	Excess minority carrier electron and hole lifetimes.

Quasi-Fermi Energy Level for Low Level Injection

For n-type Si

$$n_0 \gg p_0, \quad n_0 \gg |\delta n| = |\delta p| \gg p_0$$



$$n = n_0 + \delta n \cong n_0$$

$$p = p_0 + \delta p \cong \delta p$$

For p-type Si

$$p_0 \gg n_0, \quad p_0 \gg |\delta n| = |\delta p| \gg n_0$$



$$p = p_0 + \delta p \cong p_0$$

$$n = n_0 + \delta n \cong \delta n$$

❖ Quasi-Fermi Energy Level

$$n_0 + \delta n = N_c \exp\left(-\frac{E_c - E_{Fn}}{kT}\right) = n_i \exp\left(\frac{E_{Fn} - E_{Fi}}{kT}\right)$$

$$p_0 + \delta p = N_v \exp\left(-\frac{E_{Fp} - E_v}{kT}\right) = n_i \exp\left(\frac{E_{Fi} - E_{Fp}}{kT}\right)$$

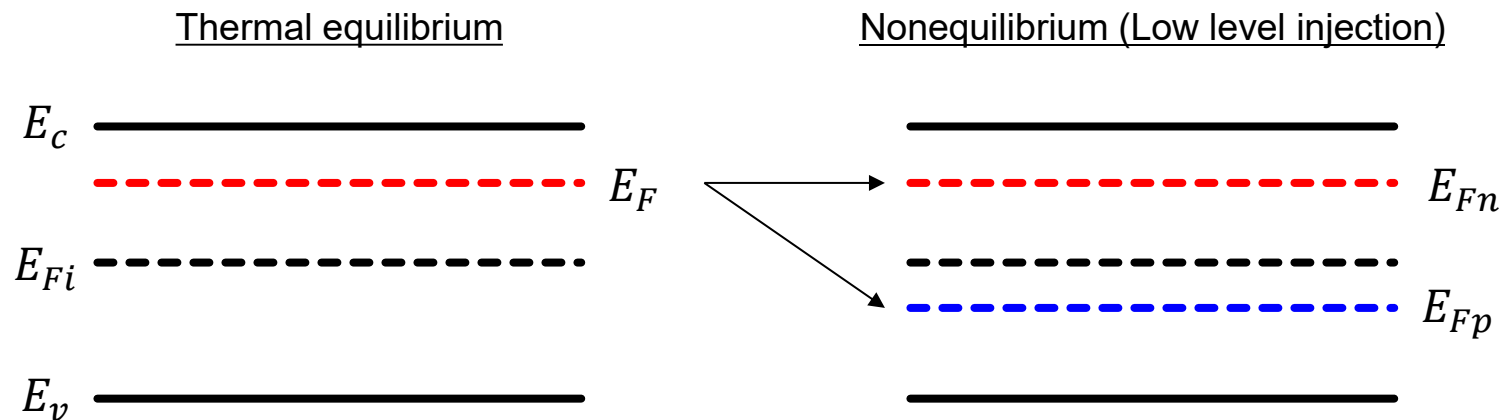
Quasi-Fermi Energy Level for Low Level Injection

❖ Quasi-Fermi Energy Level

For n-type Si

$$n_0 + \delta n = N_c \exp\left(-\frac{E_c - E_{Fn}}{kT}\right) = n_i \exp\left(\frac{E_{Fn} - E_{Fi}}{kT}\right) \cong n_0$$

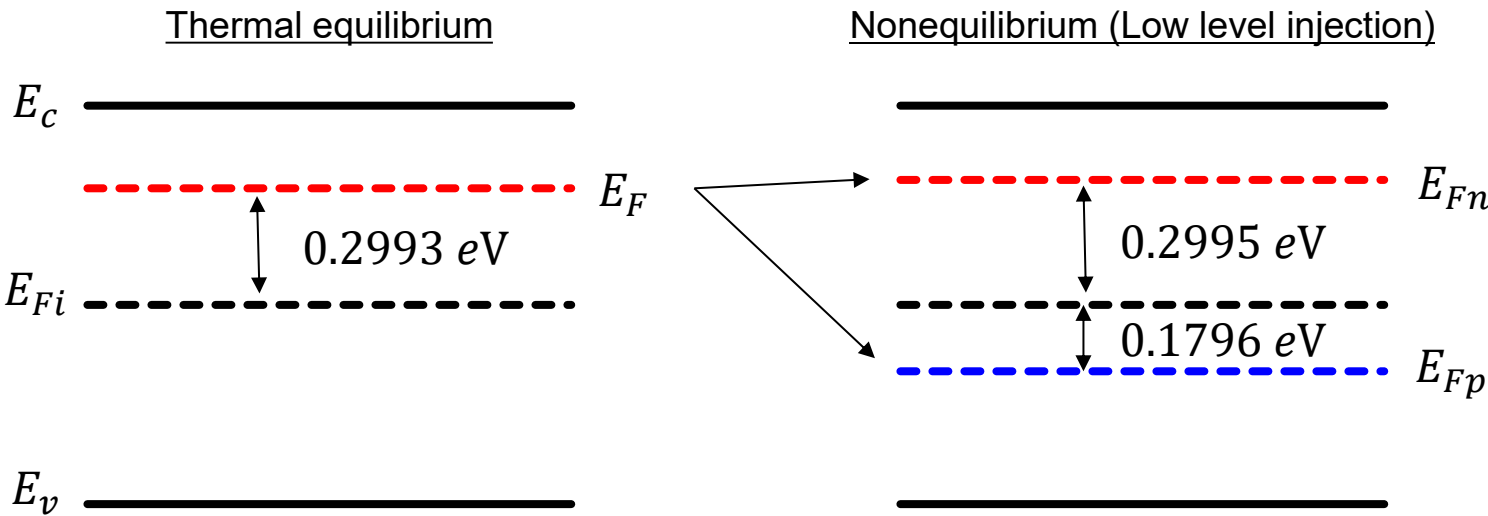
$$p_0 + \delta p = N_v \exp\left(-\frac{E_{Fp} - E_v}{kT}\right) = n_i \exp\left(\frac{E_{Fi} - E_{Fp}}{kT}\right) \cong \delta n$$



Quasi-Fermi Energy Level for Low Level Injection (Example)

Find the energy difference between the intrinsic fermi level (E_{Fi}) and the quasi-fermi energy level (E_{Fn} , E_{Fp}) for low level injection. Assume a steady-state condition.

At $T = 300\text{ K}$, $N_d = 10^{15}\text{ cm}^{-3}$, $n_i = 10^{10}\text{ cm}^{-3}$, $\tau_p = 1\text{ }\mu\text{s}$, $g'_n = 10^{19}\text{ cm}^{-3}\text{s}^{-1}$,
 $kT = 0.026\text{ eV}$, $N_c = 2.8 \times 10^{19}\text{ cm}^{-3}$, $N_v = 1.04 \times 10^{19}\text{ cm}^{-3}$



Carrier Transport Mechanisms

- Drift: Due to the Electric field
- Diffusion: Due to the Carrier concentration gradient
- Recombination-Generation (R-G): Due to the Excess carriers

Continuity Equation

Continuity Equation

→ carrier concentration change per time at the volume of $dx dy dz$

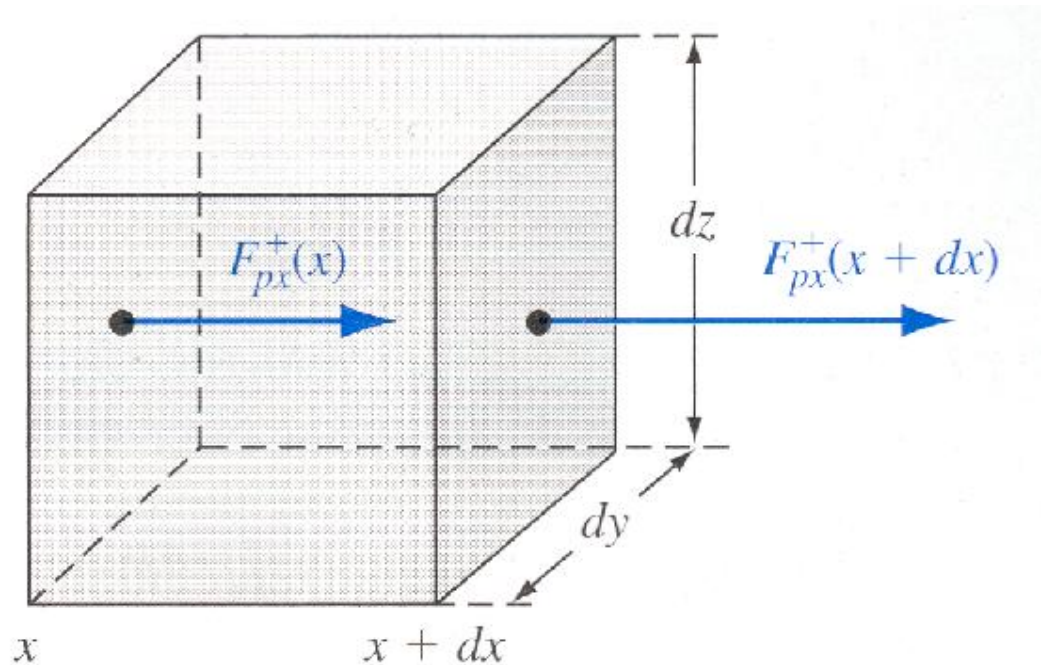


Figure 6.4 | Differential volume showing x component of the hole-particle flux.

Flux

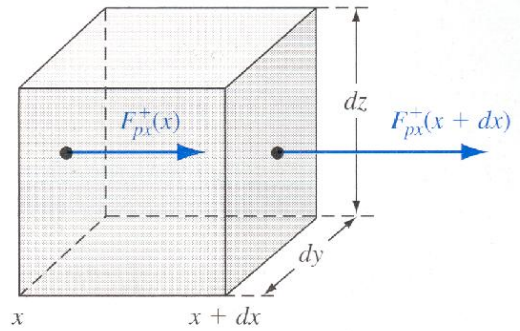
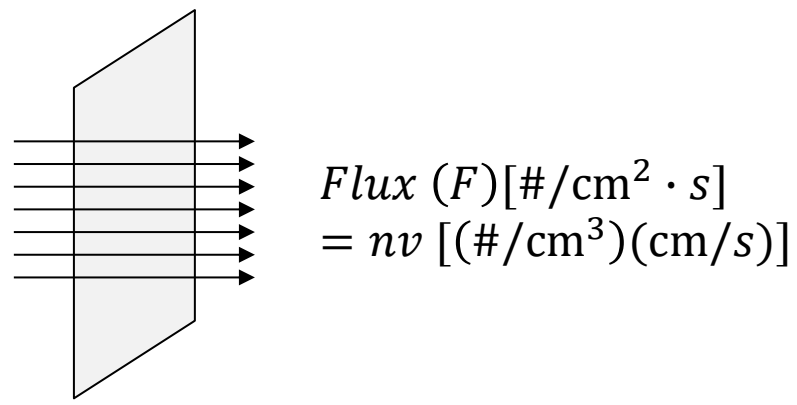


Figure 6.4 | Differential volume showing x component of the hole-particle flux.

$J = qnv = qF$

$J_n = J_{n|drf} + J_{n|diff} = qn\mu_n E + qD_n \frac{dn}{dx}$

$J_p = J_{p|drf} + J_{p|diff} = qp\mu_p E - qD_p \frac{dp}{dx}$

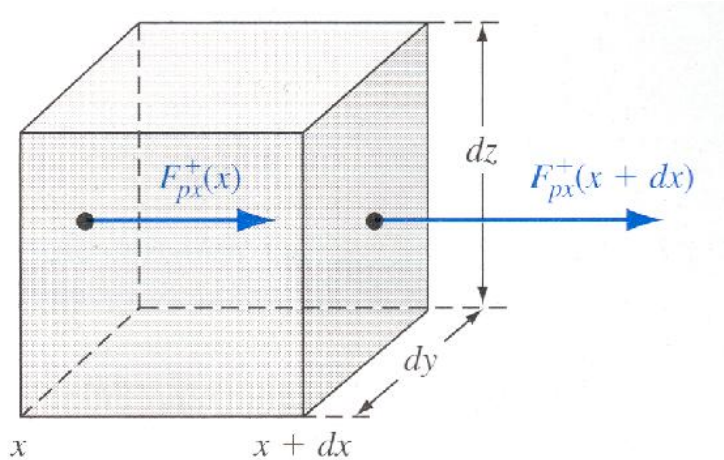
Flux

Flux

Continuity Equation

Continuity Equation

→ carrier concentration change per time at the volume of $dx dy dz$



$$\begin{aligned}\frac{\partial p}{\partial t} dx dy dz &= [F(x) - F(x + dx)] dy dz = -\frac{F(x + dx) - F(x)}{dx} dx dy dz \\ &= -\frac{\partial F(x)}{\partial x} dx dy dz\end{aligned}$$

$$\frac{\partial p}{\partial t} = -\frac{\partial F(x)}{\partial x}$$

Consider R-G

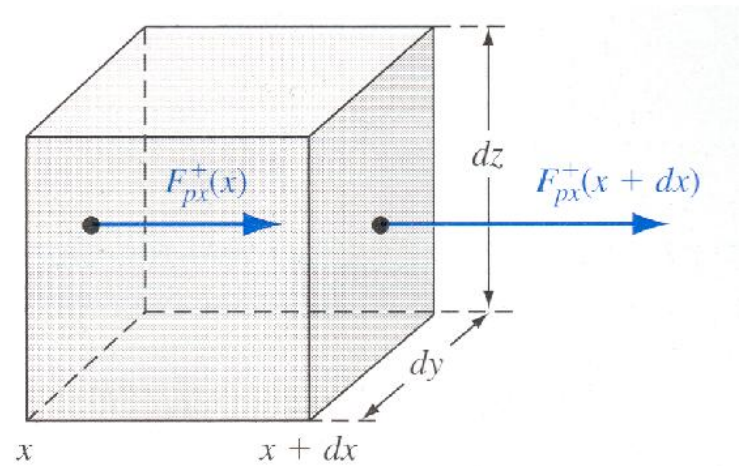


$$\frac{\partial p}{\partial t} = -\frac{\partial F(x)}{\partial x} + G_p - R_p$$

Continuity Equation

Continuity Equation

→ carrier concentration change per time at the volume of $dx dy dz$



$$\frac{\partial p}{\partial t} = - \frac{\partial F(x)}{\partial x} + G_p - R_p$$

$$J = qnv = qF(x) \qquad F(x) = \frac{1}{q}J$$

$$\frac{\partial p}{\partial t} = - \frac{\partial F(x)}{\partial x} + G_p - R_p$$

Continuity Equation

$$\frac{\partial p(t,x)}{\partial t} = - \frac{1}{q} \frac{\partial J_p}{\partial x} + g_p' - R_p' \qquad R_p' = \frac{\delta p}{\tau_p}$$

$$\frac{\partial n(t,x)}{\partial t} = + \frac{1}{q} \frac{\partial J_n}{\partial x} + g_n' - R_n' \qquad R_n' = \frac{\delta n}{\tau_n}$$



$$J_p = qp\mu_p E - qD_p \frac{dp}{dx}$$

$$J_n = qn\mu_n E + qD_n \frac{dn}{dx}$$

Continuity Equation

Continuity Equation

→ carrier concentration change per time at the volume of $dx dy dz$

$$\begin{aligned} \frac{\partial p(t,x)}{\partial t} &= -\frac{1}{q} \frac{\partial J_p}{\partial x} + g_p' - R_p' & R_p' &= \frac{\delta p}{\tau_p} \\ \frac{\partial n(t,x)}{\partial t} &= +\frac{1}{q} \frac{\partial J_n}{\partial x} + g_n' - R_n' & R_n' &= \frac{\delta n}{\tau_n} \end{aligned} \quad \Leftarrow \quad \begin{aligned} J_p(t,x) &= qp(t,x)\mu_p E(x) - qD_p \frac{\partial p(t,x)}{\partial x} \\ J_n(t,x) &= qn(t,x)\mu_n E(x) + qD_n \frac{\partial n(t,x)}{\partial x} \end{aligned}$$

$p(t,x) = p_0(x) + \delta p(t,x)$
 $n(t,x) = n_0(x) + \delta n(t,x)$

Assume
uniform doping

$p(t,x) = p_0 + \delta p(t,x)$
 $n(t,x) = n_0 + \delta n(t,x)$

For Holes

$$\frac{\partial \delta p(t,x)}{\partial t} = D_p \frac{\partial^2 \delta p(t,x)}{\partial x^2} - \mu_p \left(\frac{\partial \delta p(t,x)}{\partial x} E(x) + p(t,x) \frac{\partial E(x)}{\partial x} \right) + g_p' - \frac{\delta p(t,x)}{\tau_p}$$

For Electrons

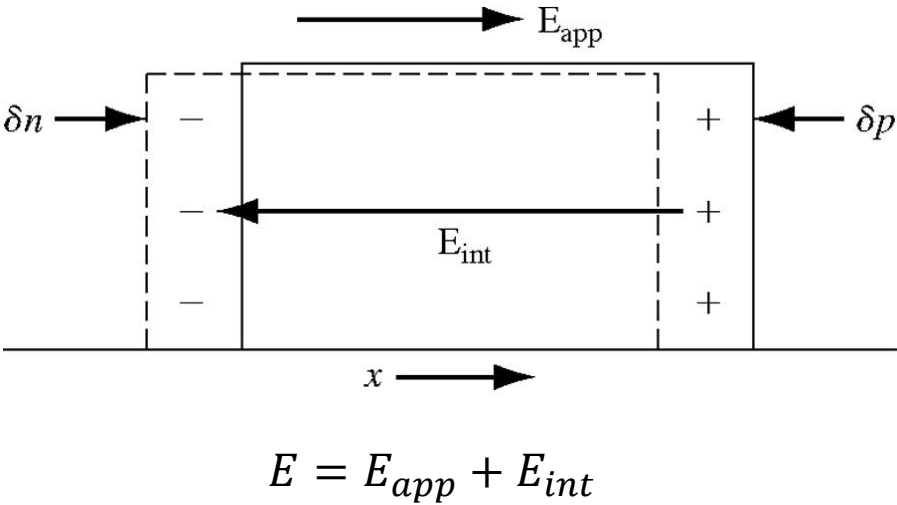
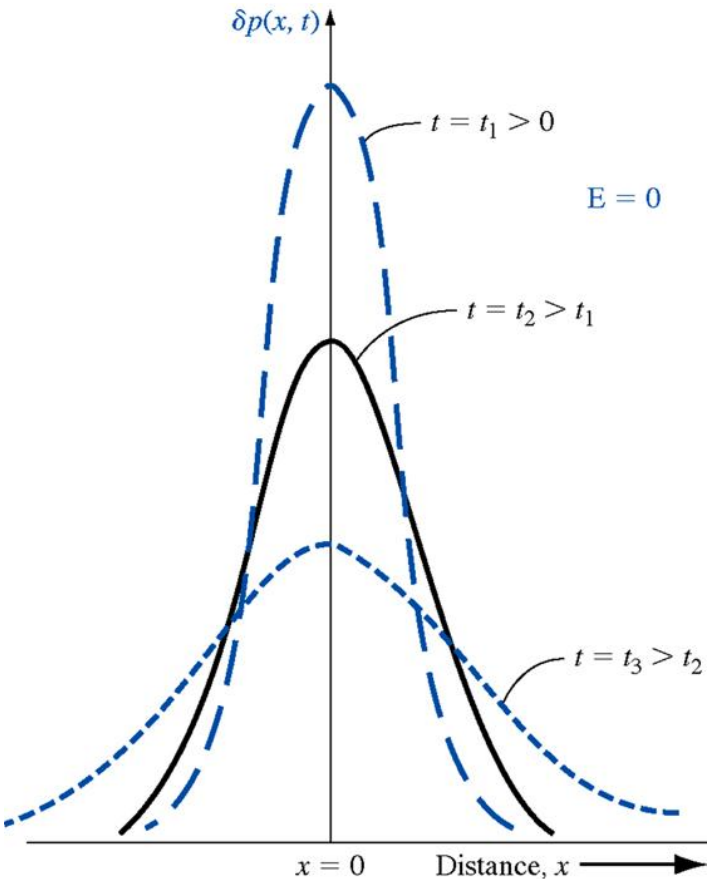
$$\frac{\partial \delta n(t,x)}{\partial t} = D_n \frac{\partial^2 \delta n(t,x)}{\partial x^2} + \mu_n \left(\frac{\partial \delta n(t,x)}{\partial x} E(x) + n(t,x) \frac{\partial E(x)}{\partial x} \right) + g_n' - \frac{\delta n(t,x)}{\tau_n}$$

Ambipolar Transport

Ambipolar transport

→ excess electron and hole move together as a pair → Single Mobility and Single Diffusion constant

 Light illumination



Ambipolar Transport

Assume

$$\delta p(t, x) = \delta n(t, x) \quad g_p' = g_n' = g' \quad \frac{\delta p(t, x)}{\tau_p} = \frac{\delta n(t, x)}{\tau_n} = R'$$

$$\frac{\partial \delta p(t, x)}{\partial t} = D_p \frac{\partial^2 \delta p(t, x)}{\partial x^2} - \mu_p \left(\frac{\partial \delta p(t, x)}{\partial x} E(x) + p(t, x) \frac{\partial E(x)}{\partial x} \right) + g_p' - \frac{\delta p(t, x)}{\tau_p}$$

× $n(t, x) \mu_n$ × $n(t, x) \mu_n$

$$\frac{\partial \delta n(t, x)}{\partial t} = D_n \frac{\partial^2 \delta n(t, x)}{\partial x^2} + \mu_n \left(\frac{\partial \delta n(t, x)}{\partial x} E(x) + n(t, x) \frac{\partial E(x)}{\partial x} \right) + g_n' - \frac{\delta n(t, x)}{\tau_n}$$

× $p(t, x) \mu_p$ × $p(t, x) \mu_p$

↓ Get rid of $\frac{\partial E(x)}{\partial x}$ $\delta N = \delta p = \delta n$

$$\begin{aligned} & (n\mu_n + p\mu_p) \frac{\partial \delta N}{\partial t} \\ &= (n\mu_n D_p + p\mu_p D_n) \frac{\partial^2 \delta N}{\partial x^2} + \mu_n \mu_p (p - n) \frac{\partial \delta N}{\partial x} E + (n\mu_n + p\mu_p) (g' - R') \end{aligned}$$

Ambipolar Transport

$$\begin{aligned} & (n\mu_n + p\mu_p) \frac{\partial \delta N}{\partial t} \\ &= (n\mu_n D_p + p\mu_p D_n) \frac{\partial^2 \delta N}{\partial x^2} + \mu_n \mu_p (p - n) \frac{\partial \delta N}{\partial x} E + (n\mu_n + p\mu_p)(g' - R') \end{aligned}$$

\Downarrow Devide by $n\mu_n + p\mu_p$

$$\frac{\partial \delta N}{\partial t} = \left(\frac{n\mu_n D_p + p\mu_p D_n}{n\mu_n + p\mu_p} \right) \frac{\partial^2 \delta N}{\partial x^2} + \left(\frac{\mu_n \mu_p (p - n)}{n\mu_n + p\mu_p} \right) \frac{\partial \delta N}{\partial x} E + g' - R'$$

\Downarrow

$$\frac{\partial \delta N}{\partial t} = \mathbf{D}' \frac{\partial^2 \delta N}{\partial x^2} + \boldsymbol{\mu}' \frac{\partial \delta N}{\partial x} E + g' - R'$$

$$\frac{\partial \delta p(t,x)}{\partial t} = D_p \frac{\partial^2 \delta p(t,x)}{\partial x^2} - \mu_p \left(\frac{\partial \delta p(t,x)}{\partial x} E(x) + p(t,x) \frac{\partial E(x)}{\partial x} \right) + g_p' - \frac{\delta p(t,x)}{\tau_p}$$

$$\frac{\partial \delta n(t,x)}{\partial t} = D_n \frac{\partial^2 \delta n(t,x)}{\partial x^2} + \mu_n \left(\frac{\partial \delta n(t,x)}{\partial x} E(x) + n(t,x) \frac{\partial E(x)}{\partial x} \right) + g_n' - \frac{\delta n(t,x)}{\tau_n}$$

$$\mathbf{D}' = \frac{n\mu_n D_p + p\mu_p D_n}{n\mu_n + p\mu_p} = \frac{D_n D_p (p + n)}{D_n n + D_p p} \text{ [cm}^2\text{/s]}$$

$$\boldsymbol{\mu}' = \frac{\mu_n \mu_p (p - n)}{n\mu_n + p\mu_p} \text{ [cm}^2\text{/V} \cdot \text{s]}$$

Einstein's Relationship

$$\frac{D_n}{\mu_n} = \frac{D_p}{\mu_p} = \frac{kT}{q} \text{ [V]}$$

Ambipolar Transport

Assume p-type Si and low-level injection

$$\begin{aligned}
 p(t, x) &= p_0 + \delta p(t, x) & p(t, x) &\cong p_0 \\
 n(t, x) &= n_0 + \delta n(t, x) & \Rightarrow n(t, x) &\cong \delta n(t, x) \\
 \delta p(t, x) &= \delta n(t, x) & p(t, x) &\gg n(t, x)
 \end{aligned}$$

For **p-type Si** and low-level injection

$$\begin{aligned}
 D' &= \frac{D_n D_p (p + n)}{D_n n + D_p p} & \text{Divide by } p D_p & \Rightarrow D' = \frac{D_n (1 + n/p)}{1 + \frac{D_n n}{D_p p}} \cong D_n
 \end{aligned}$$

$$\begin{aligned}
 \mu' &= \frac{\mu_n \mu_p (p - n)}{n \mu_n + p \mu_p} & \text{Divide by } p \mu_p & \Rightarrow \mu' = \frac{\mu_n \left(1 - \frac{n}{p}\right)}{1 + \frac{n \mu_n}{p \mu_p}} \cong \mu_n
 \end{aligned}$$

$$\frac{\partial \delta N(t, x)}{\partial t} = \textcolor{red}{D_n} \frac{\partial^2 \delta N(t, x)}{\partial x^2} + \textcolor{red}{\mu_n} \frac{\partial \delta N(t, x)}{\partial x} E + g' - R'$$

Ambipolar transport

- excess electron and hole move together as a pair → Single Mobility and Single Diffusion constant
- minority carrier determines the ambipolar transport characteristics.

Ambipolar Transport Summary

Ambipolar transport

- excess electron and hole move together as a pair → Single Mobility and Single Diffusion constant
- minority carrier determines the ambipolar transport characteristics.

For **p-type Si** and low-level injection

$$\frac{\partial \delta N(t, x)}{\partial t} = D_n \frac{\partial^2 \delta N(t, x)}{\partial x^2} + \mu_n \frac{\partial \delta N(t, x)}{\partial x} E + g' - R'$$

For **n-type Si** and low-level injection

$$\frac{\partial \delta N(t, x)}{\partial t} = D_p \frac{\partial^2 \delta N(t, x)}{\partial x^2} - \mu_p \frac{\partial \delta N(t, x)}{\partial x} E + g' - R'$$

Ambipolar Transport

Additional assumptions

- No excess carrier generation
- No external (applied) E-field
- Steady-state condition

$$g' = 0$$

$$E = 0$$

$$\frac{dn}{dt} = \frac{dp}{dt} = 0$$



Current is only generated by Diffusion and Recombination.

$$\frac{\delta N(x)}{\tau} = R'$$

For p-type Si and low-level injection

$$\frac{\partial \delta N(t, x)}{\partial t} = D_n \frac{\partial^2 \delta N(t, x)}{\partial x^2} + \mu_n \frac{\partial \delta N(t, x)}{\partial x} E + g' - R' \Rightarrow 0 = D_n \frac{d^2 \delta N(x)}{dx^2} - R'$$

For n-type Si and low-level injection

$$\frac{\partial \delta N(t, x)}{\partial t} = D_p \frac{\partial^2 \delta N(t, x)}{\partial x^2} - \mu_p \frac{\partial \delta N(t, x)}{\partial x} E + g' - R' \Rightarrow 0 = D_p \frac{d^2 \delta N(x)}{dx^2} - R'$$

Ambipolar Transport

For p-type

$$D_n \frac{d^2 \delta N(x)}{dx^2} = R'$$

$$\frac{\delta N(x)}{\tau_n} = R'$$



$$\frac{d^2 \delta N(x)}{dx^2} = \frac{\delta N(x)}{D_n \tau_n}$$

2nd order ODE

$$y'' = \frac{1}{L_N^2} y$$

For n-type

$$D_p \frac{d^2 \delta N(x)}{dx^2} = R'$$

$$\frac{d^2 \delta N(x)}{dx^2} = \frac{\delta N(x)}{D_p \tau_p}$$



$$L_n = \sqrt{D_n \tau_n} \text{ or } L_p = \sqrt{D_p \tau_p}$$

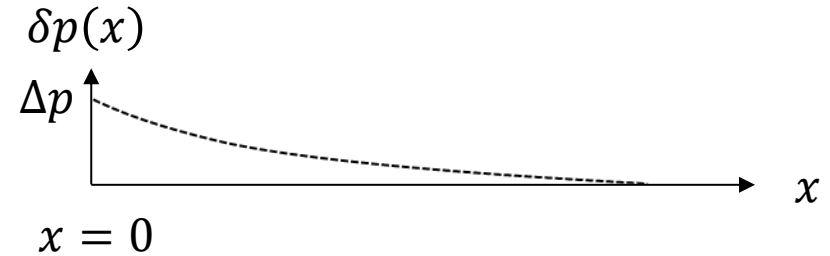
$$\delta N(x) = A e^{+\frac{x}{L_N}} + B e^{-\frac{x}{L_N}}$$

Ambipolar Transport (Example)

Additional assumptions

- No excess carrier generation
- No external (applied) E-field
- Steady-state condition

Light illumination



For n-type

$$\frac{d^2 \delta N(x)}{dx^2} = \frac{\delta N(x)}{D_p \tau} \quad L_p = \sqrt{D_p \tau} \quad \Rightarrow \quad \delta p(x) = A e^{+\frac{x}{L}} + B e^{-\frac{x}{L}}$$

Boundary condition

$$x \rightarrow \infty, \quad \delta p(x) \rightarrow 0 \quad \Rightarrow \quad A = 0$$

$$x = 0, \quad \delta p(0) = \Delta p \quad \Rightarrow \quad B = \Delta p$$

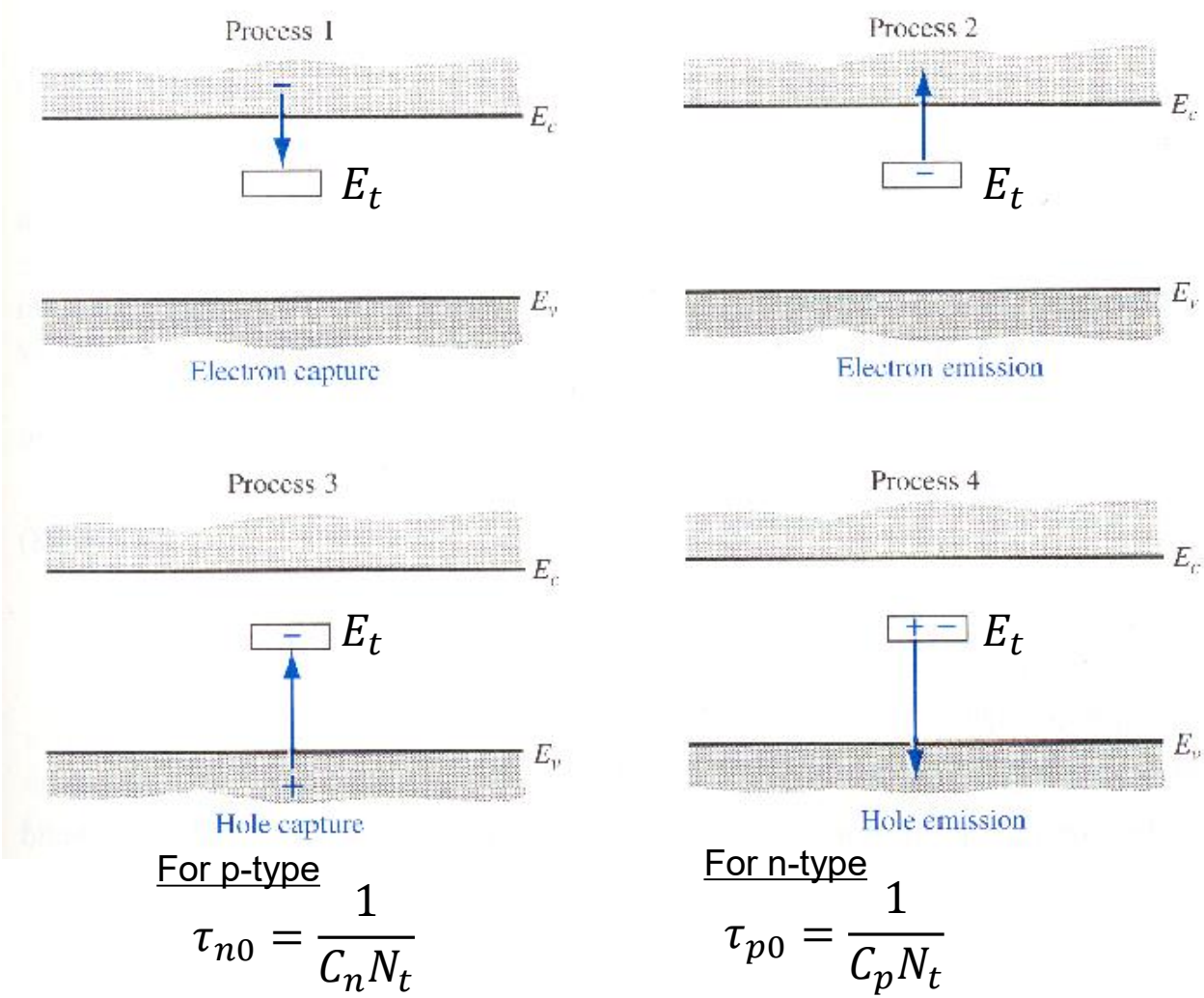
Excess carrier concentration

$$\Rightarrow \quad \delta p(x) = \Delta p e^{-\frac{x}{L_p}}$$

Current density

$$J_p(x) = J_{p|diff} = -q D_p \frac{dp(x)}{dx} = -q D_p \frac{d\delta p(x)}{dx} = \frac{q D_p \Delta p}{L_p} e^{-\frac{x}{L_p}} = q \frac{D_p}{L_p} \delta p(x)$$

Excess Carrier Lifetime



$C_n(C_p)$ = constant proportional to electron(hole) – capture cross section
 N_t = Total concnentraion of trapping centers